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More on Experimental Clarification of Maxwell-similar Gravitation Equations

Annotation

This article is a revised version of the previous article [14] to reflect the new Samokhvalov articles and other articles [20-23]. Thus, Maxwell-similar gravitation equations and the experiments of Samokhvalov are considered. It is noted that the observed effects are so significant, that in order to explain them within the said Maxwell-similar gravitation equations these equations should be supplemented by a certain empirical coefficient that may be named gravitational permeability of the medium. Then it is shown that with such supplement the results of experiments are in good agreement with so modified gravitation equations. A crude estimate of this coefficient is given. Some corollaries of these equations are considered, in particular, the gravitational excitation of electric current, the impact of gravito-magnetic induction on the electric current. Some phenomena that can be explained with the aid of these equations are indicated. This paper is a translation of paper [24].

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References

1. Introduction

There are known Maxwell equations for electromagnetic field of the form (1), proposed by Heaviside [1] (the formulas are given in Appendix 1). Heaviside is also the author the gravitation theory [2], where the gravitation field is described by equations (3) of similar form. Later it was shown [3], that in a weak gravitation field at low speeds from basic equations of general relativity he gravitational analogs of electromagnetic field may be derived, and they have the similar form (3).

Thus, in weak gravitation field of the Earth the Maxwell-similar equations may be used for the description of gravitational interactions. It means that there exist gravitational waves possessing gravito-electric component with intensity $E_g$ and gravito-magnetic component with induction $B_g$. Mass $m$, moving in a magnetic field with speed $v$, is subject to Lorentz gravito-magnetic force (analog of the known Lorentz force) of the form (in the CGS system)

$$F = \zeta \frac{m}{c} \left[ v \times B_g \right],$$

where $\zeta$ is a coefficient equal to 1 according to Heaviside, and equal to 2 in general relativity theory.

Samokhvalov [4-11, 20] had conceived and carried out a series of unexpected and surprising experiments, which presumably can be explained by interaction of irregular mass currents. Irregular mass currents $J_g$ create variable gravito-electrical intensity $E_g$ and gravito-magnetic induction $B_g$. At the interaction of this induction with the masses $m$, moving with speed $v$ there arises gravito-magnetic Lorentz force. It is important to note that the effect are so significant, that in
order to explain them within the said Maxwell-similar equations these equations should be supplemented by a certain empirical coefficient \( \xi \). Further it is shown that with such modification the results of experiments are in good agreement with the modified gravitation equations.

Thus, based on the Samokhvalov's experiments the Maxwell-similar equations should be rewritten in the form

\[
\begin{align*}
\operatorname{div}E_g &= 4\pi G m, \\
\operatorname{div}B_g &= 0, \\
\operatorname{rot}E_g &= -\frac{1}{c} \frac{\partial B_g}{\partial t}, \\
\operatorname{rot}B_g &= \frac{4\pi G \xi}{c} J_g + \frac{1}{c} \frac{\partial E_g}{\partial t}.
\end{align*}
\]

Where the value of coefficient \( \xi \) will be determined below from the said experiments. This coefficient can be called the gravitational permeability of the medium.

Lorentz force for mass

\[
F = mE_g + \xi m \frac{v \times B_g}{c},
\]

2. Certain Analogies and Consequences

Here we shall consider certain analogies between electrodynamics and gravito-electrodynamics, and some consequences of the above examined equations.

2.1. The Induction of Circular Mass Current

Magnetic flow \( \Phi \), passing through the area \( S \) of the coil with the length \( L \), carrying alternated current \( J \), in CGS system

\[
\Phi = \frac{4\pi}{c} \frac{SJ}{L}.
\]

The induction average for the area \( S \) is

\[
B = \frac{4\pi J}{cL}.
\]

If the coil is a ring of diameter \( R \), then
\[ B = \frac{2J}{cR}. \]  
(3)

Let us assume now that the ring is carrying alternated mass current \( J_g \). Then, without considering the technical realization, by analogy with (1.5) we shall get

\[ B_g = \frac{2G\xi J_g}{cR}. \]  
(4)

### 2.2. Gravitational Excitation of Electric Current

From (1.4) follows that gravitational moving force created by gravito-magnetic flow in the circuit of mass current is

\[ \xi_g = \frac{1}{c} \frac{d\Phi_g}{dt}. \]  
(5)

The force of induced electric current in a closed-loop (in the CGS system) is:

\[ J = \frac{1}{cR_{e}} \frac{d\Phi}{dt}, \]  
(5a)

where \( R_{e} \) - the resistance to these electrons motion. This current in the metal is created by free electrons with the charge \( e_o \). By analogy, taking into account (5), we find that variable gravito-magnetic flow \( \Phi_g \) also creates vortex induced mass current

\[ J_g = \frac{\xi_g}{cR_m} \frac{d\Phi_g}{dt}, \]  
(6)

where \( R_m \) is the resistance to mass particles motion. This current in the metal is created by free electrons of the mass \( m_e \). Then \( R_m = R_{e} \) - resistance to the electrons motion. In this case mass current \( J_g \) corresponds to electric current

\[ J_{ge} = J_g \frac{e_o}{m_e}. \]  
(7)

It is known that
Thus, the strength of the induced current created by variable gravito-magnetic flow $\Phi_g$ is

$$J_{ge} = \frac{\eta}{cR_e} \cdot \frac{d\Phi_g}{dt}. \tag{9}$$

Similarly to (7), the mass current $J$ corresponds to mass current

$$J_{gm} = J \frac{m_e}{e_o}. \tag{9a}$$

Thus, the strength of mass current created by variable magnetic flow $\Phi$ is

$$J_{gm} = \frac{1}{cR_e \eta} \cdot \frac{d\Phi}{dt}. \tag{9b}$$

### 2.3. Rotation of a Porous Ring

Let us consider a ring of average radius $R$, made of porous metal and charged electrically. Evidently the charges are located on the pores surfaces. Approximately we may assume that the charges distribution density along the ring's circle is described by the function:

$$\rho(\varphi) \approx \rho_o \cdot (1 + \sin(\lambda \varphi)), \tag{10}$$

where

- $\rho_o$ - is a constant,
- $\varphi$ - angular coordinate,
- $\lambda$ - the length of a "wave" depending on the average distance between the pores.

If we cause rotation of the ring with a certain angular speed $\omega$, then the charges distribution density along the circle of the ring becomes a functions of time $t$

$$\rho(t) \approx \rho_o \cdot (1 + \sin(\lambda \omega t)), \tag{11}$$

The current flowing in the ring is

$$J(t) = \frac{d\rho(t)}{dt} \approx \rho_o \cdot \lambda \omega \cdot \cos(\lambda \omega t), \tag{12}$$
where \( m_o \) is a constant. This current creates a magnetic flow that is perpendicular to the ring’s plane. The magnetic induction of this flow average by the ring’s area is determined in the CGS system by the formula (3). Consequently, the magnetic induction of a rotating charged porous ring, average by the ring’s area, is

\[ B \approx 2\rho_o \omega \lambda \cdot \cos(\lambda \omega t)/(cR). \] (13)

By analogy we can state that the rotating porous ring creates mass current

\[ J_g(t) = \frac{dm(t)}{dt} \approx m_o \cdot \lambda \omega \cdot \cos(\lambda \omega t). \] (14)

Then from (4) we find that this current creates variable gravito-magnetic induction

\[ B_g \approx 2m_o \xi G \omega \lambda \cdot \cos(\lambda \omega t)/(cR). \] (15)

### 2.4. Induction of a Moving Body

It is known that the induction of field in a medium with permeability \( \mu \), created by a charge \( q \), moving with speed \( v \), in a certain point, is

\[ \overline{B} = \mu q(v \times r)/cr^3. \] (16)

The vector \( r \) is directed from the point, where the moving charge \( q_1 \) is located, to the referred point. Similarly, the gravito-magnetic induction of the field created by the mass \( m \), moving with a speed \( \bar{v} \), in a certain point, is

\[ \overline{B_g} = \xi G m(\bar{v} \times \bar{r})/cr^3, \] (17)

Because, as shown in the Section 2.2, the electronic current is at the same time also the mass current, the gravito-magnetic induction can create the Lorentz force, affecting the electric current.

### 2.5. Gravitomagnetic law the Biot-Savart-Laplace

It is known that an electric current creates a magnetic flux density, determined by the Biot-Savart-Laplace law as

\[ \overline{dB} = \frac{\mu \cdot J}{r^3 c} \left[ dL \times \overline{r} \right] \] (18a)

where \( \overline{dL} \) - vector element conductor with current, \( \overline{r} \) - vector between it and the point where, it is determined of induction. This law is currently
being considered as a consequence of Maxwell's equations. Therefore, it can be argued that a similar law for gravitomagnetic induction, generated of mass current. In this case, the Biot-Savart-Laplace law is written as follows:

$$\overline{dB_g} = \frac{\xi Gm}{r^3 c} \left[ \vec{v} \times \vec{r} \right], \quad (18a)$$

where $\vec{v}$ - the speed vector of the mass $m$.

### 2.6. Gravito-magnetic Ampere Force

It is known that a conductor carrying electric current $\vec{J}$ in a magnetic field with induction $\vec{B}$ is affected by Ampere force (per a length unit

$$\overline{F_a} = \frac{1}{c} \left( \vec{J} \times \vec{B} \right) \quad (19)$$

Similarly, a conductor carrying mass current $\vec{J}_g$ in a gravito-magnetic field with induction $\vec{B}_g$ is affected by Ampere force

$$\overline{F_{ag}} = \frac{\xi}{c} \left[ \vec{J}_g \times \vec{B}_g \right], \quad (20)$$

Let us consider the case when mass current is a consequence of electric current, i.e. the particles carrying the charge form the mass current. Then

$$\vec{J}_g = J \eta_2, \quad (21)$$

$$\eta_2 = \frac{m}{q}, \quad (22)$$

where $m, q$ – mass and charge of the particle. Then a conductor carrying electric current $\vec{J}$ in a gravito-magnetic field with induction $\vec{B}_g$ is affected by Ampere force

$$\overline{F_{age}} = \frac{\xi \eta_2}{c} \left[ \vec{J} \times \vec{B}_g \right]. \quad (23)$$

For example, if the charged particle is an electron, then

$$m_e \approx 9.1 \cdot 10^{-34} \text{ g}, \quad e_o \approx 1.6 \cdot 10^{-19} \text{ Kl},$$

$$\eta_2 = \frac{m_e}{e_o} \approx 0.6 \cdot 10^{-14} \frac{\text{ g}}{\text{ Kl}}. \quad (24)$$
But if the charged particle is an ion with mass \( m = h \cdot m_e \), then
\[
\eta_2 = \frac{h \cdot m_e}{e_0} \approx 0.6h \cdot 10^{-14} \frac{g}{\text{Kl}}.
\] (25)
and for complex molecules \( \eta_2 \gg 1 \). So, at the interaction of gravito-magnetic induction with electrical current significant Ampere forces are likely to act.

2.7. Density of Magnetic Wave Energy

It is known that the density of electromagnetic wave energy [13], is
\[
W = \frac{B^2}{8\pi} \left[ \frac{g}{\text{sm} \cdot \text{sec}^2} \right].
\] (26)
By applying the derivation shown there for the equations (1.2-1.5) of gravito-electromagnetic wave, we find
\[
W_g = \frac{B_g^2}{8\pi G}.
\] (27)

2.8. Induction of Current-carrying Conductor

It is known that the magnetic induction of infinite conductor carrying electric current is:
\[
B = 2J / (cd),
\] (28)
where \( d \) - is the distance from the conductor to the point of measurement. Similarly, the gravito-magnetic induction of infinite conductor with mass current is
\[
B_g = 2\xi G J_g / (cd).
\] (29)

3. Certain Experimental Estimates

The analysis of Samokhvalov's experiments [4-11, 20], performed in Appendix 2, permits to obtain a crude estimate of the coefficient \( \xi \) of gravitational permeability. There it was shown that for vacuum
\[
\xi \approx 10^{10}.
\] (30)
This value can be greatly understated, as the experiments were carried out at an average vacuum, but \( \xi \) increases with decreasing pressure. For atmospheric pressure \( \xi \gg 0 \), which explains the absence of visible effects of gravitational interaction of moving masses.
There exist several phenomena which can be explained with the aid of the equations (1.2-1.5), considered above, with the existence of coefficient of gravitational permeability $\xi$ – see [12, 15-19, 21-23].

Appendixes
In each application of the formula numbered independently, and links to these formulas are written as "(p.'application number and partition it')"

Appendix 1. The Equations of Electro-magnetism and Gravito-magnetism
Further we shall use the following notations:

- $q$ - electric charge $[g \cdot sm]$;
- $\rho$ - electric charge density $[g \cdot sm^3]$;
- $J$ - electric current density $[1 \cdot [g \cdot sm] / [sm \cdot sec]]$;
- $c$ - speed of light in vacuum $c \approx 3 \cdot 10^10 [sm/sec]$;
- $E$ - electric field intensity $[v \cdot sm / sec^2 = 3 \cdot 10^4 V/m]$;
- $B$ - magnetic induction $[1 / [sec \cdot sm] = G_s]$;
- $\varepsilon$ - permittivity of the medium is equal to 1 for the vacuum in the CGS system;
- $\mu$ - permeability of the medium is equal to 1 for the vacuum in the CGS system;
- $v$ - speed $[sm/sec]$;
- $F$ - force $[dyn = g \cdot sm/ sec^2]$;
- $m$ - mass $[g]$;
- $\rho_g$ - mass density $[g/sm^3]$;
- $J_g$ - mass current density $[g/sm^2/sec]$;
- $G$ - gravitational constant,

$G \approx 7 \cdot 10^{-8} \left[ \frac{dyn \cdot sm^2}{g^2} \right] = \frac{sm^3}{g \cdot sec^2}$;
The Maxwell equations for electromagnetism in CGS system are as follows [1]:

\[
\begin{align*}
\text{div} E &= 4\pi \rho, \\
\text{div} B &= 0, \\
\text{rot} E &= -\frac{1}{c} \frac{\partial B}{\partial t}, \\
\text{rot} B &= \frac{4\pi \cdot \mu}{c} J + \frac{\varepsilon}{c} \frac{\partial E}{\partial t}.
\end{align*}
\]

The Lorentz force for the electric charge is

\[
F = qE + \frac{q}{c} [v \times B].
\]

The Maxwell equations for gravito-electromagnetism in medium in Gauss CGS system [3], supplemented by analogy with equations (1-4) permeability \( \xi \), are as follows:

\[
\begin{align*}
\text{div} E_g &= 4\pi G \rho_g, \\
\text{div} B_g &= 0, \\
\text{rot} E_g &= -\frac{1}{c} \frac{\partial B_g}{\partial t}, \\
\text{rot} B_g &= \frac{4\pi G \xi}{c} J_g + \frac{1}{c} \frac{\partial E_g}{\partial t}.
\end{align*}
\]

The Lorentz force for the mass is

\[
F = mE_g + \xi \frac{m}{c} [v \times B_g],
\]

where \( \xi \) - is a coefficient equal to 1 by Heaviside and equal to 2 in general relativity theory.

**Appendix 2. The Experiments of Samokhvalov**

**1. Experiment 1**

We shall consider the experiment of Samokhvalov described in [4]. Two disks are placed into a vacuum chamber; they are misbalanced (by
skewed axes) and are rotating in one direction. Both disks are overheated. Technical parameters of the setup are as follows:

- Material of the disks: aluminum
- Pressure in the chamber: 1 Pa
- Density of aluminum: \( \rho \approx 2.7 \text{ g/sm}^3 \)
- Thickness of the disks: \( h \approx 0.09 \text{ sm} \)
- Diameter of the disks: \( 2R = 16.5 \text{ sm} \)
- Gap between the disks: \( d \approx 0.3 \text{ sm} \)
- Beating on the sides: \( 0.05 \text{ sm} \)
- Number of revolutions: \( f \approx 50/\text{sec} \)
- Temperature of overheating (in [4] is written that the temperature rise measured after some minutes was 50 K).

Let us consider the disk's rotation as mass current. We can assume that this current is formed by the mass's motion in the circle of the upper band of the disk of radius \( R \approx 7 \text{ sm} \) and the cross-section

\[
S \approx 0.3 \cdot 2.5 \text{sm}^2 \approx 7.5 \text{sm}^2.
\]  

(1)

The speed of this mass is

\[
v = 2\pi R \cdot f \approx 2\pi \cdot 7 \cdot 50 \approx 2200 \text{sm/sec}.
\]  

(2)

So, the mass current is

\[
J_g = S \rho v \approx 7.5 \cdot 2.7 \cdot 2200 = 4400 \text{g/sec}.
\]  

(3)

This current is variable because the beating of the disks. In accordance with (2.4) this current causes a variable axial induction (along the \( ox \) axis of the disk) average on the circle area of radius \( R \),

\[
B_g = \frac{2\xi G J_g}{cR}
\]  

(4)

or

\[
B_g = \frac{2 \cdot \xi \cdot 7 \cdot 10^{-8} \cdot 4400}{3 \cdot 10^{10} \cdot 7} \approx 3 \xi \cdot 10^{-15}.
\]  

(5)

This induction is variable in time because of the disks. We shall assume that the circular frequency of this induction is

\[
\omega \approx 2\pi f = 314.
\]  

(6)

In accordance with (2.9), the strength of vortex electric current created by variable gravito-magnetic flow, is
\[ J_{ge} = \frac{\eta}{cR_e} \cdot \frac{d\Phi_g}{dt} \]  

(7)

or

\[ J_{ge} = \frac{\eta \omega}{cR_e} \cdot \Phi_g . \]  

(8)

In our case

\[ \Phi_g = \beta \pi R^2 B_g = \beta \pi R^2 \cdot 3 \cdot 10^{-15} , \]  

(9)

where \( \beta \) – is the coefficient of induction weakening on the level of the driven disk (because of the gap). So,

\[ J_{ge} = \frac{\eta \omega}{cR_e} \cdot \beta \pi R^2 B_g \]  

(10)

or

\[ J_{ge} = \frac{1.8 \cdot 10^{14} \cdot 314}{3 \cdot 10^{10} R_e} \cdot \beta \pi 8.25^2 \cdot 3 \cdot \xi 10^{-15} = \frac{\xi \beta}{R_e} 10^{-6} \]  

(10a)

This electric current raises the disk temperature. In the experiment it was shown that the disk's temperature has increased by \( \Delta T \approx 100 \) grades. Let us consider the equivalent voltage

\[ E_e = J_{ge} R_e \]  

(11)

And assume that such increase of the disk temperature may be due to the voltage \( E_e \). From (10a, 11) we find

\[ E_e = \xi \beta 10^{-6} . \]  

(12)

Let us assume that such equivalent voltage is \( E_e = 200 \). Then we find

\[ \xi \beta \approx 2 \cdot 10^8 . \]  

(13)

Here \( \xi \) depends on the pressure, and \( \beta \) depends on the gap. Assuming that \( \beta \approx 1/d^2 \) and knowing that \( d \approx 0.3sm \), we find \( \beta \approx 0.01 \). Thus, based on Samokhvalov's experiment we can now assume that for the indicated conditions the gravitational permeability coefficient with the pressure of 0.1 atm is equal to

\[ \xi_p(0.1) \approx 2 \cdot 10^{10} . \]  

(14)
2. Experiment 2

Let us now consider the experiments of Samokhvalov described in [5]. Two disks are placed into a vacuum chamber, misbalanced by skewed axes. The first of them rotates forcibly, and the second disk begins rotation due to the impact of the first one. The speed \( f_2 \) of the second disk's rotation (if the rotation speed of the first one is constant) depends on the gap between the disks \( d \) and on the pressure in vacuum chamber \( p \). We may assume that the rotation speed of the driven disk is

\[
f_2(p,d) = f_2(p) \cdot f_2(d).
\]

This experiment explores these two dependences.

The dependence of rotation speed on the pressure is given in [5] on Fig 2, from which we find

\[
p= [0.1, 0.3, 0.5, 0.7, 0.9, 1] \text{ (atm)},
\]

\[
f= [24, 17, 8, 2, 0.2, \varepsilon],
\]

where \( \varepsilon \) is a small value that it is impossible to find from the experiment results.
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Fig. 1 shows this experimental dependence (by circles) and (by full line) – the approximating function in the form of a polynomial with 5 members. We assume that,
\[
f_2(p, d = 0.2) = f_{2p}(p) \cdot f_{2d}(0.2)
\]
(2)

In particular, by approximating function we find:
\[
f_2(0.1, 0.2) = 25, \quad f_2(0, 0.2) \approx 35.
\]
(2a)

The dependence of rotation speed on the distance

is given in [5, Fig. 3], from which we find:
\[
\begin{align*}
d &= [0.15, 0.2, 0.25, 0.3] \text{ (sm)}, \\
f_1 &= [24, 17, 6, 5] \text{ при } p = 1 \text{ atm}, \\
f_{102} &= [30, 25, 12, 10] \text{ при } p = 1.02 \text{ atm}.
\end{align*}
\]

Fig. 2 shows this experimental dependence (by circles) and the approximating function (by full line) – in the form of \( a + b/d^2 \), and the function
\[
f_{2d}(d) = 1/d^2.
\]
(3)

To a first approximation further we shall use the function (2). In particular, for \( d = 0.2 \) (cm) we have
\[
f_{2d}(0.2) \approx 25.
\]
(3a)
Analysis of the functions $f_{2p}(p)$ and $f_{2d}(d)$

Taking into account (2, 3), we find:

$$f_{2p}(p) = f_{2}(p, 0.2)/f_{2d}(0.2) = 0.04 f_{2}(p, 0.2).$$  \hspace{1cm} (4)

In particular from (2a) we find:

$$f_{2p}(0) = 0.04 f_{2}(0, 0.2) = 0.04 \cdot 35 \approx 1.5,$$  \hspace{1cm} (6)

Below in (p.3.7) it will be shown that

$$f_{2p}(p) = g \cdot \xi_{p}^{2}(p).$$  \hspace{1cm} (8)

Thus,

$$\xi_{p}(p) \approx \sqrt{\frac{f_{2p}(p)}{\theta}},$$  \hspace{1cm} (9)

From (9) it follows that

$$\frac{\xi_{p}(0)}{\xi_{p}(p)} \approx \sqrt{\frac{f_{2p}(0)}{f_{2p}(p)}},$$  \hspace{1cm} (10)

In experiment 1 it was shown, that

$$\xi_{p}(0.1) \approx 2 \cdot 10^{10}.$$

Combining (5, 6, 10, 11), we get

$$\xi_{p}(0) \approx \xi_{p}(0.1) \sqrt{\frac{f_{2p}(0)}{f_{2p}(0.1)}} \approx 2 \cdot 10^{10} \sqrt{\frac{1.5}{1}} \approx 2.5 \cdot 10^{10}.$$

From this we can find a crude estimate of the gravitational permeability of vacuum:

$$\xi \approx 10^{10}.$$  \hspace{1cm} (13)

3. The Role of Gravito-magnetic Lorentz Forces

In Samokhvalov's experiments the driving disk drags the driven disk. Now we shall present the explanation of this phenomenon. Samokhvalov notes that first there occurs the vibration of the driving disk, and then begins the rotation of the driven disk – then see Fig. 3.

The disks' vibration is explained in the following way - see Fig. 3. Above, analyzing the Experiment 1, it was shown that the driving disk is a variable mass current (p.2.1.3) with circular frequency (p.2.1.6). This current mass $m_{i}$, moving with speed $V_{1}$, creates a variable gravito-magnetic induction (p.2.1.4), which is perpendicular to the mass current...
of drive disc, ie radially and parallel to the disc plane – see a closed curve on Fig. 3. This induction vector at the slave drive moves with a speed $v_1$ relative to the mass $m_2$ of the driven disc. This raises gravito-magnetic Lorentz force, acting on the mass $m_2$ and directed vertically and having the form

$$ F_1 = m_2 v_1 B_g \frac{\xi}{c}. $$

(1)

![Fig. 3.](image)

Above, when analyzing the experiment 1, we have showed that the masses $m_1, m_2$ are the mass of a circle of higher band of the disk with radius $R \approx 7 \text{cm}$ and cross-section (p.2.1.1). This mass is equal to
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\[ m_1 = m_2 = 2\pi R S \rho. \]  \hspace{1cm} (2)

The force \( F_1 \) is directed perpendicularly to the disk plane and varies with the frequency \( f \approx 50/sec \), causing the vibration of the driven disk. Evidently, the speed \( v_2 \) of this vibration is proportional to the force \( F_1 \), i.e.

\[ v_2 = \alpha F_1, \]  \hspace{1cm} (3)

where \( \alpha \) is a certain constant.

This force may explain the "oscillatory" character of the process of repulsion of the screen with the increase of the oscillations amplitude (angle of the frame's deviation) after steadying of the disk rotation speed", which is reflected in the Samokhvalov's experiments described in \([8]\).

**Rotating force acting on the driven disk** is explained as follows – see Fig. 3. The foregoing gravito-magnetic induction \( B_g \) (p.2.1.4), created by the driving disk is directed perpendicularly to the mass current of the driving disk, i.e. along the disk's radius and parallel to its plane. This induction acts on the vertically vibrating mass \( m_2 \) of the driven disk by gravito-magnetic Lorentz force (2.18):

\[ F_2 = m_2 v_2 B_g \frac{\xi}{c}. \]  \hspace{1cm} (4)

This force is tangential to the circumference of the disc, because perpendicular to the direction of induction (which is directed along the radius of the disk) and the speed (which is perpendicular to the plane of the disk). Due to the fact, that the speed of vibration \( v_2 \) and the induction \( B_g \) are changing synchronously, the vector of this force doesn't change direction. Apparently, the rotation speed of the driven disk is proportional to the force \( F_2 \), i.e. the number of its revolutions is

\[ f_2 = \gamma F_2, \]  \hspace{1cm} (5)

where \( \gamma \) – a certain constant. Combining (1-5) we get

\[ f_2 = \gamma m_2 v_2 B_g \frac{\xi}{c} = \gamma m_2 B_g \frac{\xi}{c} \alpha F_1 = \]  

\[ = \gamma m_2 B_g \frac{\xi}{c} \alpha m_2 v_1 B_g \frac{\xi}{c} = \alpha \gamma \left( m_2 \frac{\xi}{c} B_g \right)^2. \]  \hspace{1cm} (6)
Because gravito-magnetic induction $B_g$ proportional gravito-magnetic permeability $\xi$ (which follows from (p.1.9) in Appendix 1), the number of revolutions of the slave drive is

$$f_2 = g \cdot \xi^2.$$  (7)

Which is proportional to $\xi^2$ with a certain proportionality factor $g$. This ratio is used in the above analysis of the Experiment 2 – see (p.2.2.8).

### Appendix 3. Some formulas in the CGS system

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<tr>
<th>Name</th>
<th>Electromagnetism</th>
<th>Gravitomagnetism</th>
</tr>
</thead>
<tbody>
<tr>
<td>Maxwell's equations</td>
<td>div$E = 4\pi \rho / \varepsilon$</td>
<td>div$E_g = 4\pi G \rho_g$</td>
</tr>
<tr>
<td></td>
<td>div$B = 0$</td>
<td>div$B_g = 0$</td>
</tr>
<tr>
<td></td>
<td>rot$E = -\frac{1}{c} \frac{\partial B}{\partial t}$</td>
<td>rot$E_g = -\frac{1}{c} \frac{\partial B_g}{\partial t}$</td>
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<tr>
<td></td>
<td>rot$B = \left( \frac{4\pi \cdot \mu}{c} J \right)$</td>
<td>rot$B_g = \left( \frac{4\pi G \xi}{c} J_g \right)$</td>
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<tr>
<td>Lorentz force</td>
<td>$F = qE + \frac{q}{c} [v \times B]$</td>
<td>$F = mE_g + \xi \frac{m}{c} [v \times B_g]$</td>
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<tr>
<td>Magnetic flow passing through the area of coil with a current (p.2.1)</td>
<td>$\Phi = \frac{4\pi \mu}{c} \cdot \frac{SJ}{L}$</td>
<td>$\Phi_g = \frac{4\pi G \xi}{c} \cdot \frac{SJ_g}{L}$</td>
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<tr>
<td>Induction of the ring current (p.2.1)</td>
<td>$B = \frac{2\mu J}{cR}$</td>
<td>$B_g = \frac{2G \xi J_g}{cR}$</td>
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<tr>
<td>The moving force (p.2.2)</td>
<td>$\varepsilon = \frac{1}{c} \cdot \frac{d\Phi}{dt}$</td>
<td>$\varepsilon_g = \frac{1}{c} \cdot \frac{d\Phi_g}{dt}$</td>
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<tr>
<td>The strength of the induced current (p.2.2)</td>
<td>$J = \frac{1}{cR_e} \cdot \frac{d\Phi}{dt}$</td>
<td>$J_g = \frac{1}{cR_m} \cdot \frac{d\Phi_g}{dt}$</td>
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</tbody>
</table>
### Induction of a moving body (p.2.4)

\[ B = \mu q (v \times r) / cr^3 \]

\[ B_g = \xi Gm (v \times r) / cr^3 \]

### Biot-Savart-Laplace law (p.2.5)

\[ dB = \frac{\mu \cdot J}{r^3 c} [dL \times r] \]

\[ dB_g = \frac{\xi Gm}{r^3 c} [v \times r] \]

### Ampere Force (p.2.6)

\[ F_a = \frac{1}{c} (J \times B) \]

\[ F_{ag} = \frac{1}{c} [J_g \times B_g] \]

### The energy density of a magnetic wave (p.2.7)

\[ W = \frac{B^2}{8\pi} \]

\[ W_g = \frac{B_{g}^2}{8\pi G} \]

### Induction of current conductor (p.2.8)

\[ B = 2J / (cd) \]

\[ B_g = 2\xi GJ_g / (cd) \]

### References

17. Khmelnik S.I. Active field honeycomb, ibid.

24. Khmelnik S.I. More on Experimental Clarification of Maxwell-similar Gravitation Equations, ibid